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Article in *Proceedings of SPIE - The International Society for Optical Engineering* · November 2003

DOI: 10.1117/12.524039

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Finite object Talbot effect as a lens produced image

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ABSTRACT

In this work we theoretically calculate the self-image field amplitude of a finite periodic object. It is compared with the field of the image of an unitary cell of the same object formed by a lens. The results are verified by simulations of the two processes.

1. INTRODUCTION

When we illuminate an infinite periodic object of transmittance $t(x,y)$ with a monochromatic plane wave, at a distance Z_N in front of the object, we find a plane where the field amplitude is identical to the transmittance function of the object [1]. These planes are called ‘‘ Talbot’’ planes or ‘‘self-images’’. There also exist other planes Z_{NT} between them where the field amplitude reproduces $t(x,y)$ but shifted half period (See Figure 1).

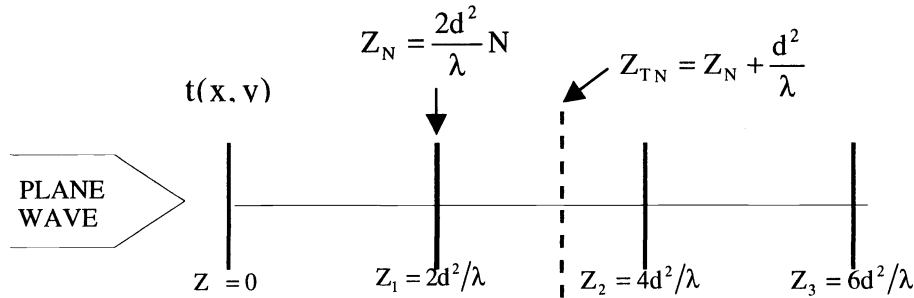


Figure 1. Talbot Effect.

Where d is the period, N is a natural number and λ is the wave length of the incident light.

2. THEORETICAL DESCRIPTION

The transmittance function of an ideal sampling filter, i.e. an infinite periodic object formed by points that are separated a distance d , can be represented as [2,3]:

$$T_1(x, y) = \left[\exp\left(-i \frac{k}{2f} (x^2 + y^2)\right) \otimes t_u(x, y) \right] \cdot A(x + nd, y + md) \quad (1)$$

Where (n,m) represents the point coordinates (row, column), f is the self-image distance in Talbot effect, \otimes is the convolution operation and $\kappa = 2\pi/\lambda$. Now we consider the transmittance of a ‘‘finite’’ periodic object with unitary cells $t_u(x, y)$ with pupil function $A(x, y)$. For the cell (n,m) , in particular, we obtain the transmittance function:

$$T_1(x, y) = \left[\exp\left(-i \frac{k}{2f} (x^2 + y^2)\right) \otimes t_u(x, y) \right] \cdot A(x + nd, y + md) \quad (2)$$

Using the Fresnel approximation, the properties of the Fourier transform and of the convolution, we obtain, for the object whose transmittance is given by equation (2), a field amplitude in the self-image plane:

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$$U(x, y) = A \cdot \exp\left(\frac{iK}{2f}(x^2 + y^2)\right) \cdot \left\{ \left(t_u(x, y) \cdot \exp\left(\frac{-iKx^2}{2f}\right) \right) \otimes \left(T_A\left(\frac{x}{\lambda f}\right) \cdot \exp\left(\frac{i2\pi d}{\lambda f}(nx + my)\right) \right) \right\} \quad (3)$$

Where A is a constant and $T_A(x/\lambda f)$ is the Fourier transform of the periodic object pupil function.

3. IMAGE FORMED BY A LENS

If we illuminate an unitary cell with a spherical wave and we use a lens to form its image (See figure 2), the field amplitude in the image plane is identically to the expression of the equation (3).

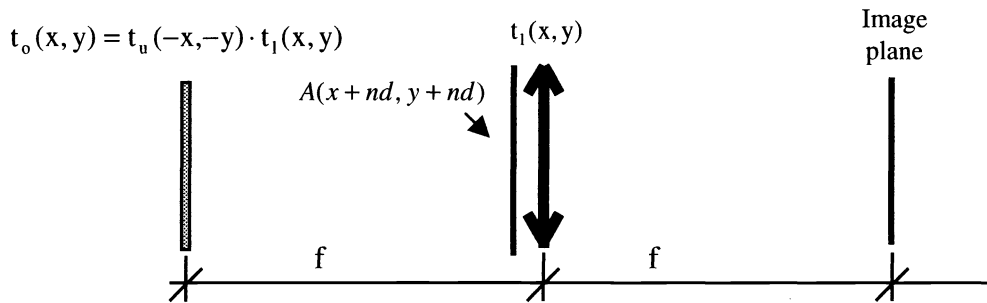


Figure 2. Unitary cell image by a lens

Where $t_l(x, y)$ is the lens transmittance whose focal length is $f/2$ and $A(x+nd, y+md)$ is the lens pupil function equivalent to the finite size object in Talbot effect.

4. SIMULATIONS

Making the simulation of the processes that are shown in figures 1 and 2 we can observe the coincidence among the field amplitude for an unitary cell in Talbot effect (Figure 3) with the image of an unitary cell formed by a lens (Figure 4).

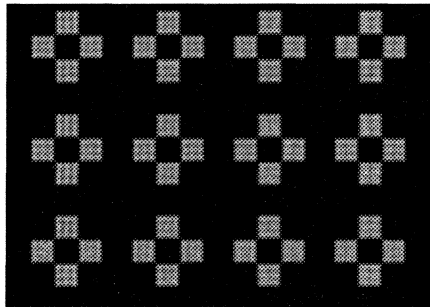


Figure 3

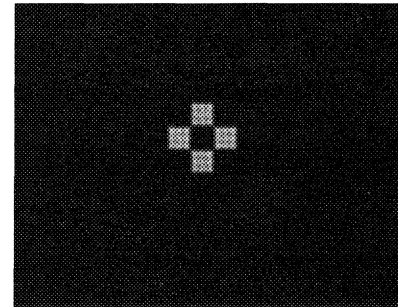


Figure 4

5. CONCLUSIONS

Equation (3) shows that the field amplitude in the self-image plane is affected by the finite size of the object, since the quality of the image falls toward the borders when having larger values for n and m. We tested, by using simulations, (Figures 3 and 4) that the unitary cell field amplitudes in both cases are the same.

6. REFERENCIAS

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